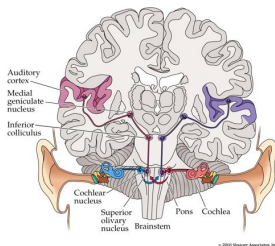


We do interferometry all the time. . .

Ascending auditory pathways



“Significant number of nerve fibers cross the brain and make connections with neurons on the side opposite from the side of the ear in which they begin. This happens very early on in the auditory system. Inter-aural comparisons are an important source of information for the auditory system about where a sound came from.”

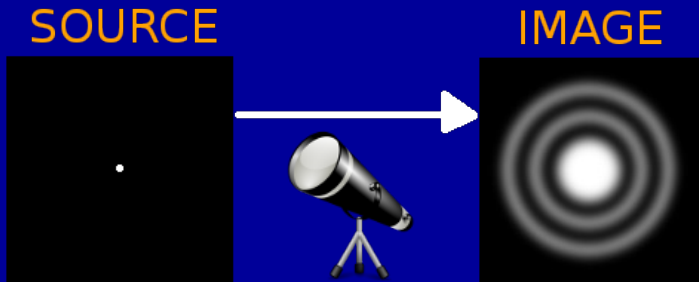
David Heeger's lecture from

<http://www.cns.nyu.edu/~david/courses/perception/lecturenotes/localization/localization.html>

Resolution of an optical device

Any optical device is fundamentally limited by **diffraction**

Due to diffraction, the image of a point source is not a point image, but a blurry point-like shape called **Point Spread Function (PSF)**



The PSF of an optical device

The PSF of an instrument is equal to the Fourier transform of the aperture's autocorrelation (in units of λ)

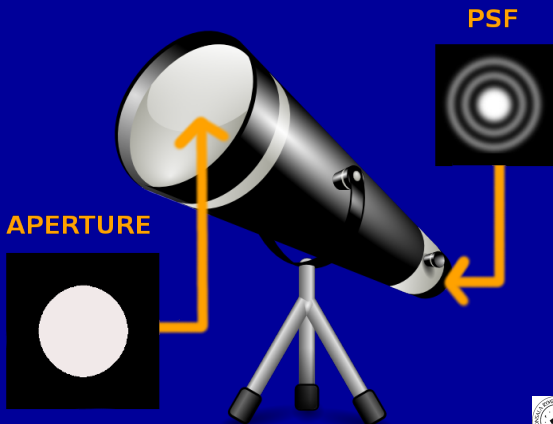
$$\text{PSF} = \mathcal{F}(\text{APERTURE} * \text{APERTURE})$$



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The PSF of an instrument is equal to the Fourier transform of the aperture's autocorrelation (in units of λ)

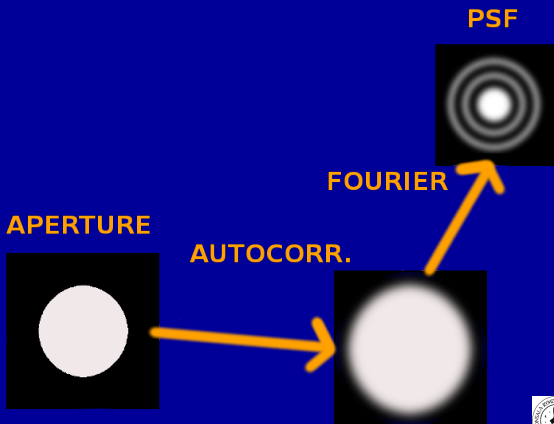
$$\text{PSF} = \mathcal{F}(\text{APERTURE} * \text{APERTURE})$$



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The PSF of an instrument is equal to the Fourier transform of the aperture's autocorrelation (in units of λ)

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The PSF of an optical device

The size of the PSF depends on the wavelength and the aperture size:

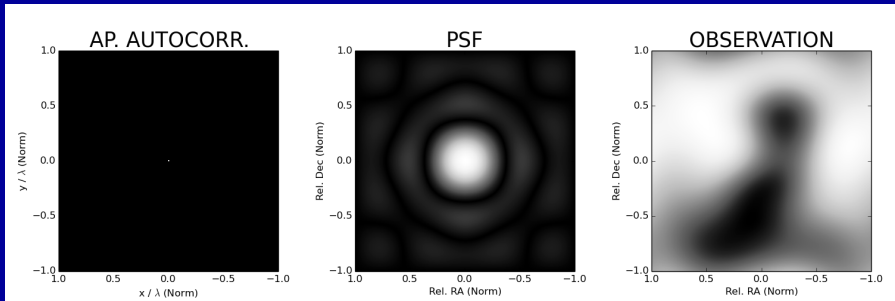
$$\Delta\theta \sim \frac{\lambda}{D}$$



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$$\Delta\theta \sim \frac{\lambda}{D}$$

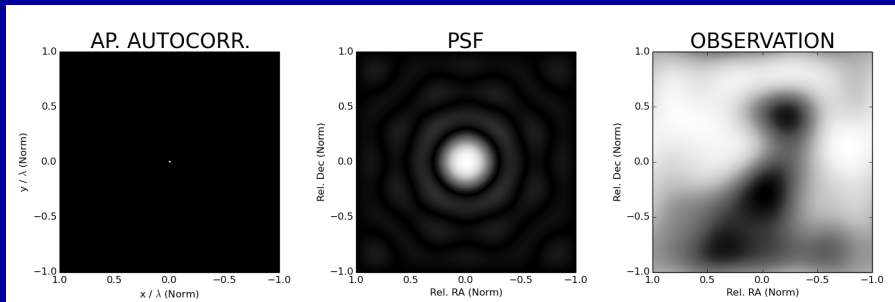


The resolution increases if the aperture increases (in units of λ)

The PSF of an optical device

The size of the PSF depends on the wavelength and the aperture size:

$$\Delta\theta \sim \frac{\lambda}{D}$$



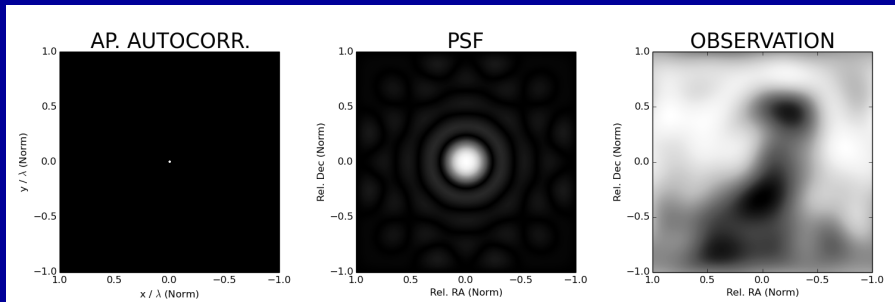
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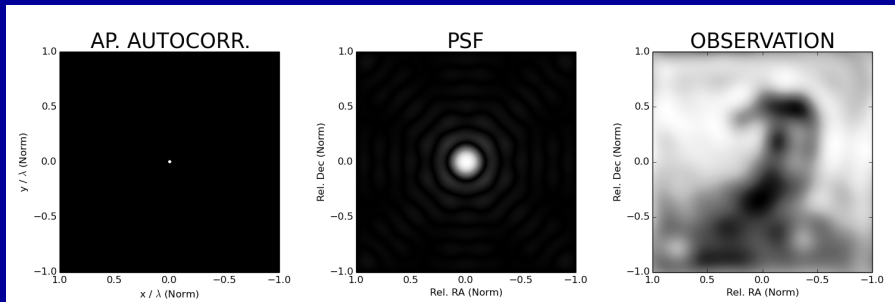
The resolution increases if the aperture increases (in units of λ)



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$$\Delta\theta \sim \frac{\lambda}{D}$$



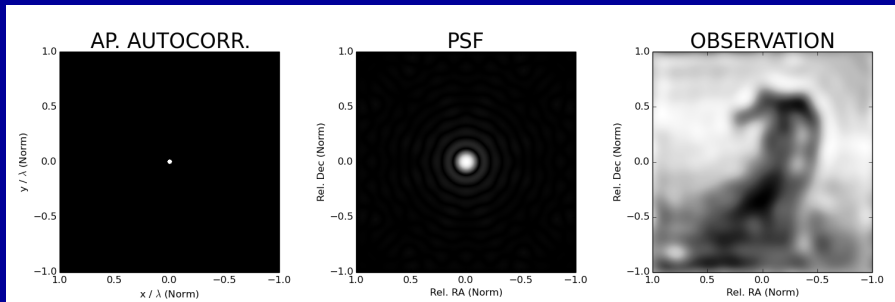
The resolution increases if the aperture increases (in units of λ)



The PSF of an optical device

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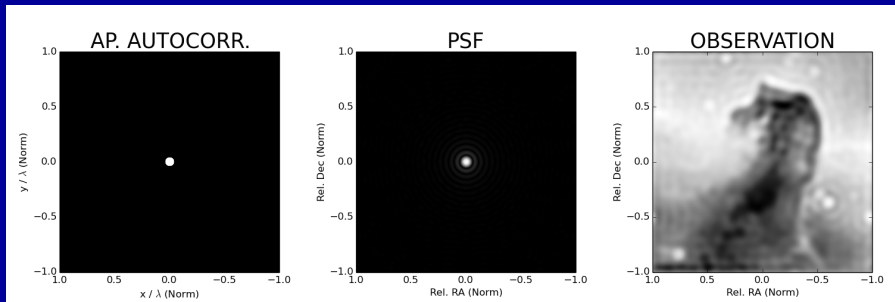
The resolution increases if the aperture increases (in units of λ)



The PSF of an optical device

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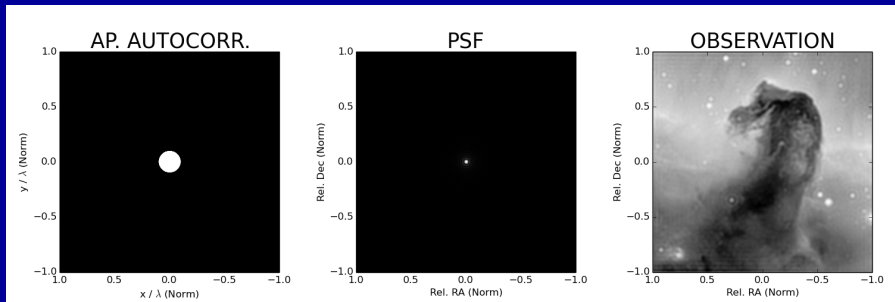
The resolution increases if the aperture increases (in units of λ)



The PSF of an optical device

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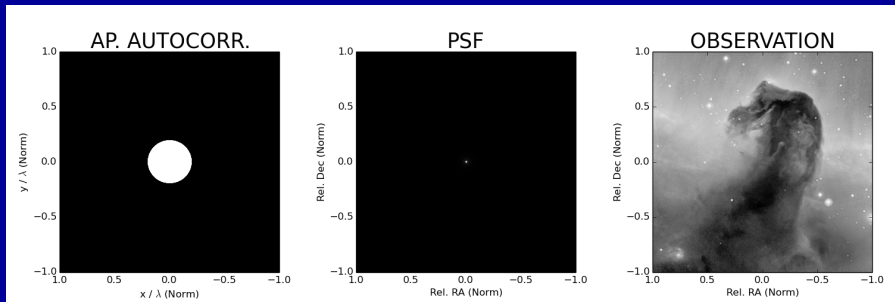
The resolution increases if the aperture increases (in units of λ)



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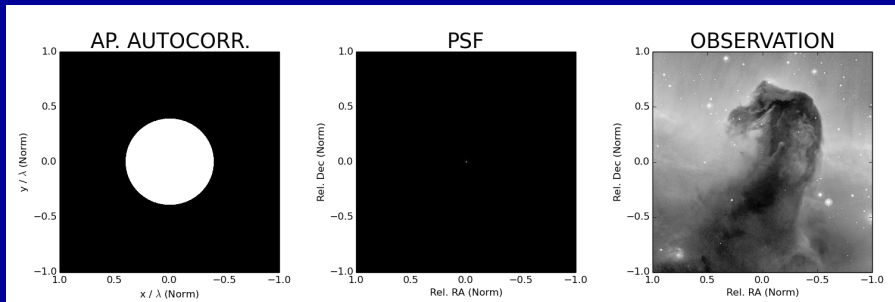


The resolution increases if the aperture increases (in units of λ)

The PSF of an optical device

The size of the PSF depends on the wavelength and the aperture size:

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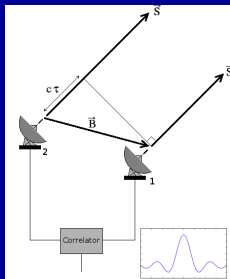
The resolution increases if the aperture increases (in units of λ)



The two-element interferometer

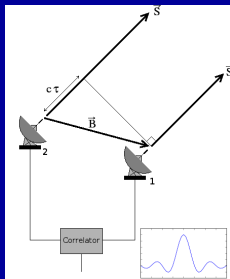


Observing with a two-element interferometer



A BASIC OBSERVABLE OF AN INTERFEROMETER IS THE
CROSS-SPECTRUM (a.k.a. **COMPLEX VISIBILITY**):

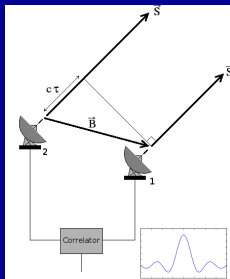
Observing with a two-element interferometer



A BASIC OBSERVABLE OF AN INTERFEROMETER IS THE
CROSS-SPECTRUM (a.k.a. **COMPLEX VISIBILITY**):

$$E_i(t)$$

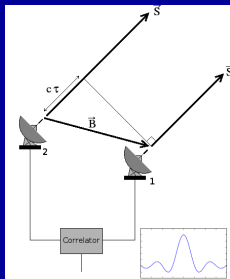
Observing with a two-element interferometer



A BASIC OBSERVABLE OF AN INTERFEROMETER IS THE
CROSS-SPECTRUM (a.k.a. **COMPLEX VISIBILITY**):

$$E_i(t) \rightarrow$$

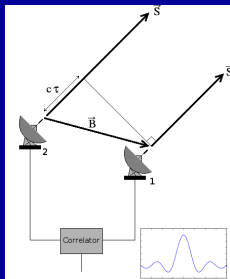
Observing with a two-element interferometer



A BASIC OBSERVABLE OF AN INTERFEROMETER IS THE
CROSS-SPECTRUM (a.k.a. **COMPLEX VISIBILITY**):

$$E_i(t) \rightarrow S_i(\nu) = \mathcal{F}(E_i(t))$$

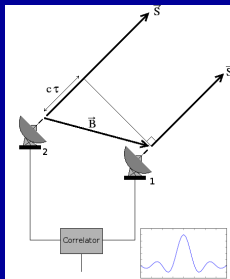
Observing with a two-element interferometer



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$$E_i(t) \rightarrow S_i(\nu) = \mathcal{F}(E_i(t)) \rightarrow$$

Observing with a two-element interferometer



A BASIC OBSERVABLE OF AN INTERFEROMETER IS THE **CROSS-SPECTRUM** (a.k.a. **COMPLEX VISIBILITY**):

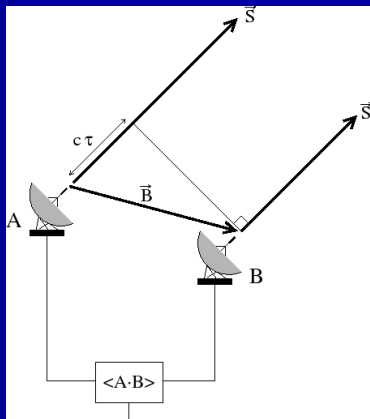
$$E_i(t) \rightarrow S_i(\nu) = \mathcal{F}(E_i(t)) \rightarrow V_{ik}(\nu) = S_i(\nu) \times S_k^*(\nu)$$

Notice the similarity with the **power spectrum** used in **single dish**:

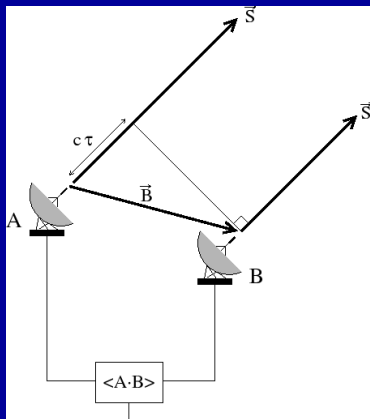
$$P_i(\nu) = S_i(\nu) \times S_i^*(\nu) = |S_i(\nu)|^2$$



Cross-spectrum of a monochromatic point source



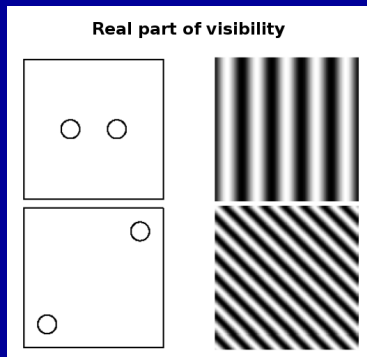
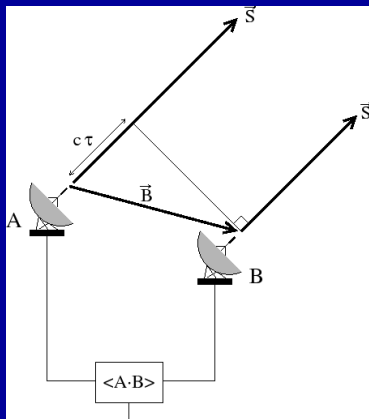
Cross-spectrum of a monochromatic point source



$$\tau = \frac{\vec{B} \cdot \vec{S}}{c}$$



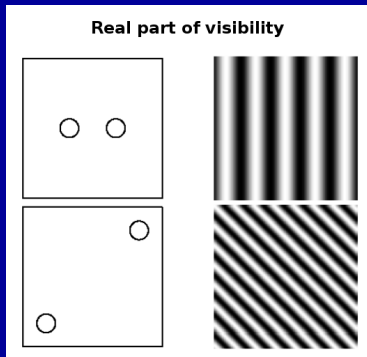
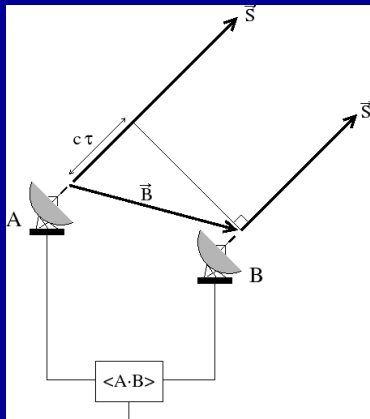
Cross-spectrum of a monochromatic point source



$$\tau = \frac{\vec{B} \cdot \vec{S}}{c} \rightarrow V(\nu_0) = I \times (\cos 2\pi\nu_0\tau - j \sin 2\pi\nu_0\tau)$$



Cross-spectrum of a monochromatic point source



$$\tau = \frac{\vec{B} \cdot \vec{S}}{c} \rightarrow V(\nu_0) = I \times (\cos 2\pi\nu_0\tau - j \sin 2\pi\nu_0\tau) = I \times \exp(-2\pi j\nu_0\tau)$$



Visibilities of extended sources

For extended sources, the overall measured visibility is the integral of the source structure times the response.

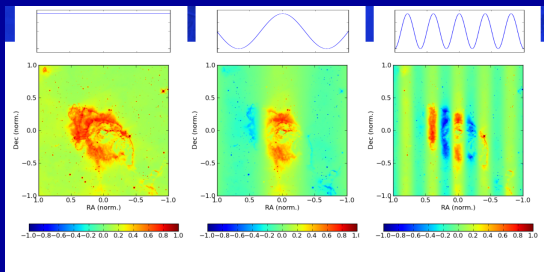
$$V(\vec{B}, \nu_0) = \int I_{\nu_0}(\vec{S}) \exp\left(-2\pi j \frac{\vec{B}\vec{S}}{\lambda_0}\right) d\vec{S}$$



Visibilities of extended sources

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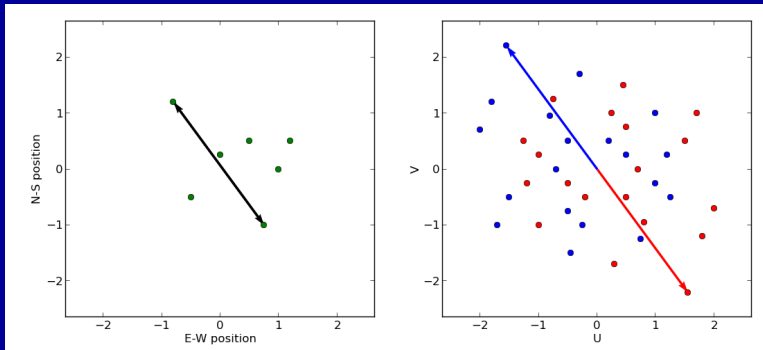
For small fields of view, this is just the integral of the source structure multiplied by a *complex fringe*.

The two-element interferometer. Summary

- **Diffraction limit** is one of the main limitations of optical devices. It is specially important at long wavelengths, where **extremely large apertures** are needed for a reasonable angular resolution.
- The **cross-spectrum** between the signals of two telescopes (observing the same source) is directly related to the Fourier transform of the source structure, computed at the point of **Fourier space** given by the **baseline projection** of the telescopes into the **UV plane** (i.e., the plane perpendicular to the source direction).
- This result will allow us to **synthesize large apertures** through **interferometric observations** (see next lecture).



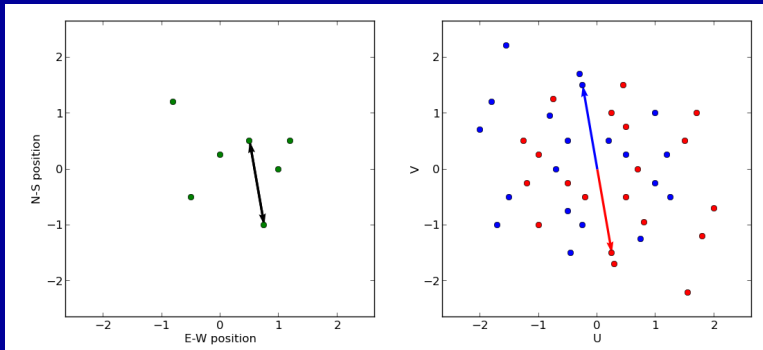
The UV plane



Each pair of telescopes is separated by its **baseline** vector. The measured visibility is the Fourier transform of the source structure measured at the points of the UV plane given by the baseline. **A multi-element interferometer is, hence, just a set of many two-element interferometers!**



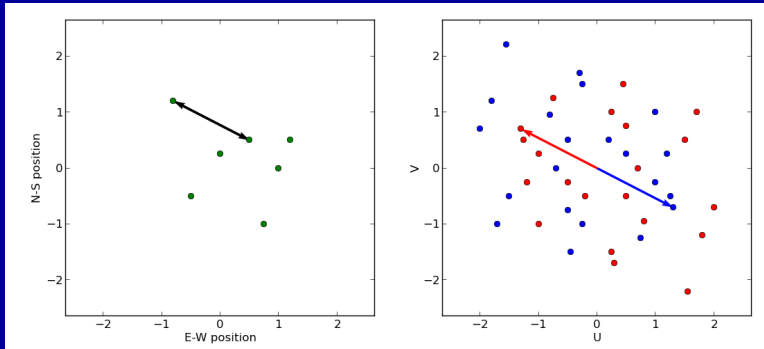
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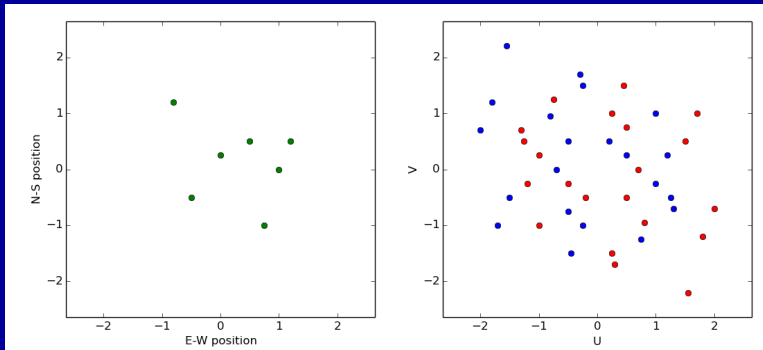


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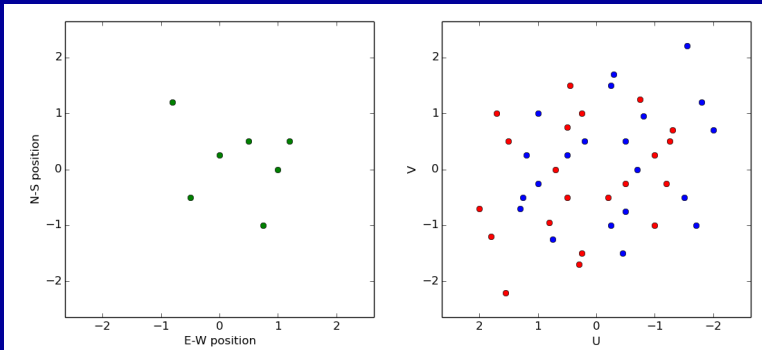
The UV plane



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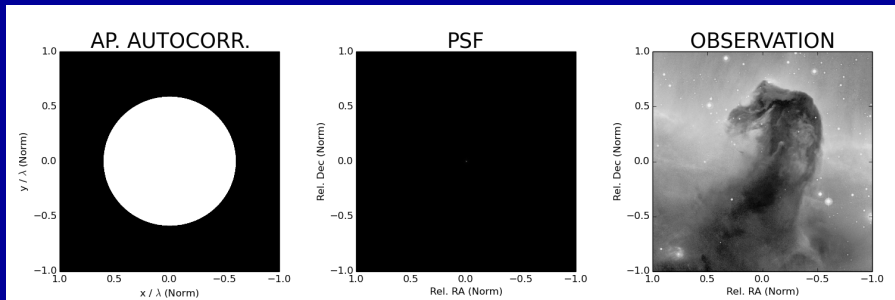
BEWARE! The **U** axis must point to **East** (like the **Right Ascension**!)



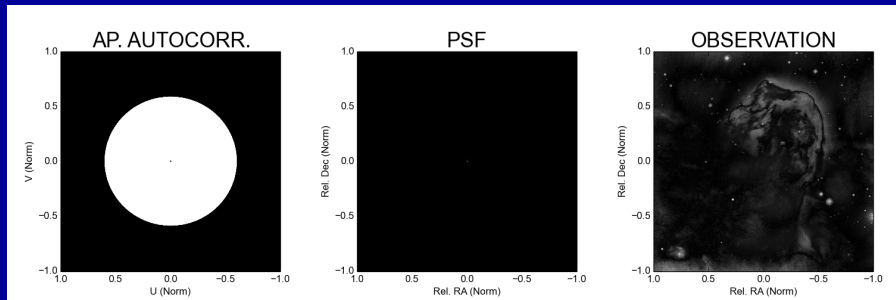
Synthesizing aperture.



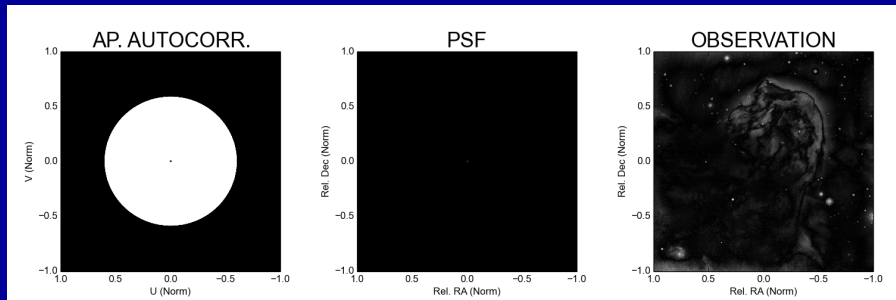
The zero-spacing hole



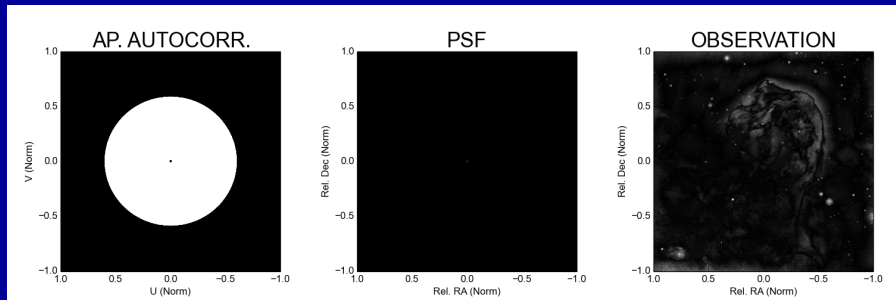
The zero-spacing hole



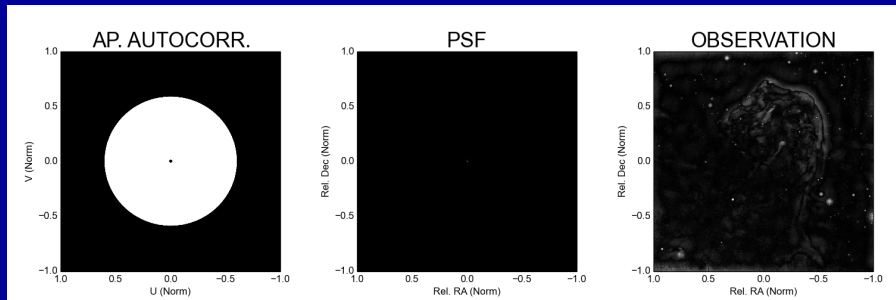
The zero-spacing hole



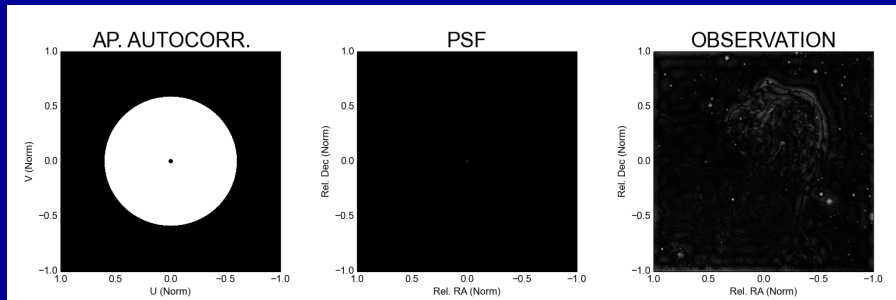
The zero-spacing hole



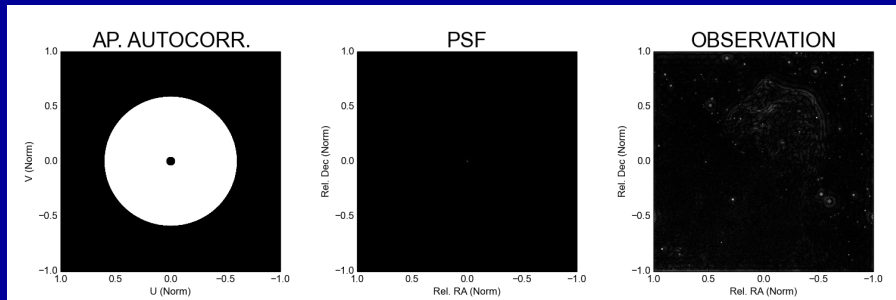
The zero-spacing hole



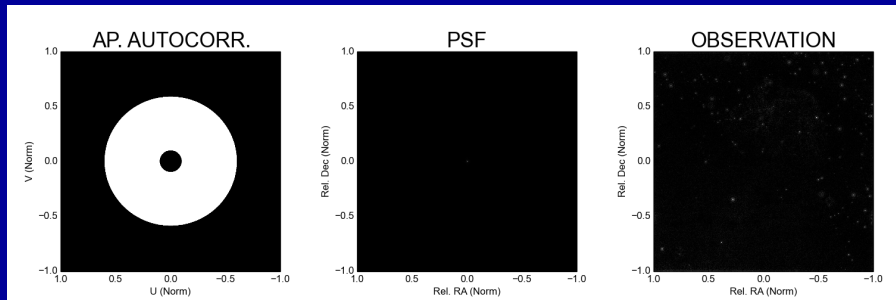
The zero-spacing hole



The zero-spacing hole



The zero-spacing hole

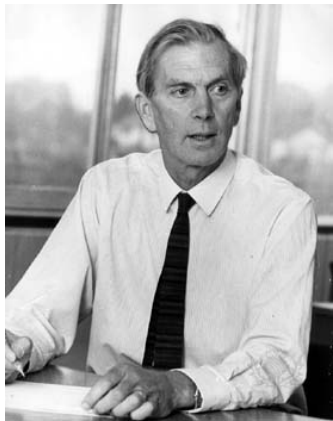


Summary

- An interferometer made of **many elements** (antennas) can be understood as a set of **many two-element** interferometers (as many as antenna pairs can be formed).
- The different antenna pairs sample **different points** of the UV plane.
- The **Earth rotation** changes the baseline projections into the source plane, hence **synthesizing** a more complete sampling of the UV plane (UV coverage).
- There are, however, **remaining holes** in the UV plane with no observations, and this will limit the quality and fidelity of our images (dynamic-range limit).
- Special care must be taken with the hole at the center of the UV plane, since it limits our sensitivity to extended structures.



How it all begun (here and there)



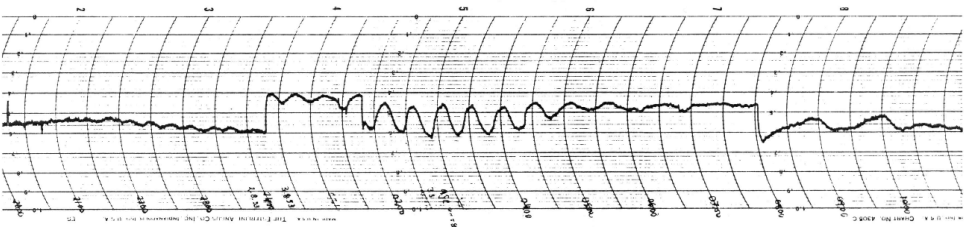
Sir Martin Ryle (1918 - 1984)
(G3CY)



Jorma J. Riihimaa (1933 - 2011)
(OH2PX, OH8PX)

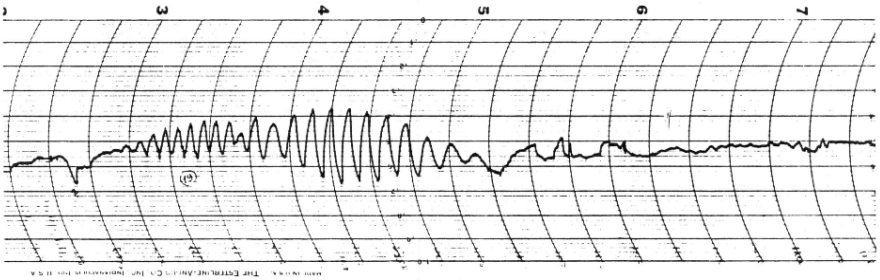
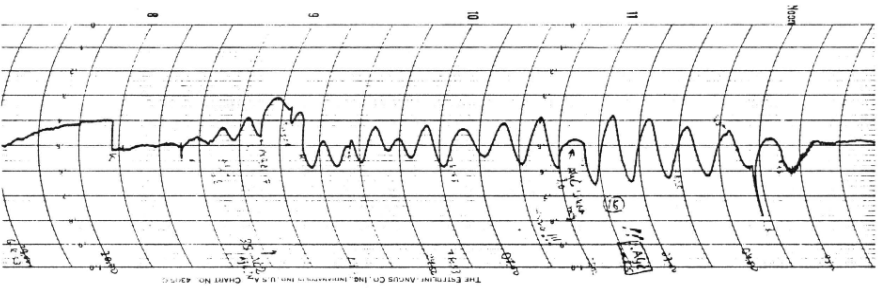


K.39. Teekkari J.J. Riihimaa etuvahvistinta virittämässä.



K.43. Ensimmäinen rekisteröinti Cass A:n ohikulusta
2-3.8.1953. (81.5 MHz)

1953



K.44. Kojeliston parantelu. Oikealla näkyvät
 Cyg A:n ja Cass A:n perättäiset ohikulut
 (pienemmällä rek.pap. nopeudella).

